# Imaging Sequences I

#### M219 - Principles and Applications of MRI Kyung Sung, Ph.D. 2/19/2025

# **Course Overview**

- 2025 course schedule
  - https://mrrl.ucla.edu/pages/m219\_2025
- Assignments
  - Homework #3 due on 3/5

- TA office hours, Mon 4-6pm
- Office hours, Fri 10-11am

RF Pulse Bandwidth and Slice Profile: Small Tip Angle Approximation

## Bloch Equation (at on-resonance)

$$\frac{d\vec{M}_{rot}}{dt} = \vec{M}_{rot} \times \gamma \vec{B}_{eff}$$
where  $\vec{B}_{eff} = \begin{pmatrix} B_1(t) \\ 0 \\ B_0 - \frac{\omega}{\gamma} + G_z z \end{pmatrix}$ 

When we simplify the cross product,

$$\frac{d\vec{M}}{dt} = \begin{pmatrix} 0 & \omega(z) & 0\\ -\omega(z) & 0 & \omega_1(t)\\ 0 & -\omega_1(t) & 0 \end{pmatrix} \vec{M}$$
$$\omega(z) = \gamma G_z z \quad \omega_1(t) = \gamma B_1(t)$$

$$\begin{aligned} & \frac{d\vec{M}}{dt} = \begin{pmatrix} 0 & \omega(z) & 0 \\ -\omega(z) & 0 & \omega_1(t) \\ 0 & -\omega_1(t) & 0 \end{pmatrix} \vec{M} \\ & M_z \approx M_0 \text{ small tip-angle approximation} \\ & \sin \theta \approx \theta \\ & \cos \theta \approx 1 \\ & M_z \approx M_0 \rightarrow \text{constant} \end{aligned} \right\} \quad \frac{dM_z}{dt} = 0 \\ & \frac{M_{xy}}{dt} = -i\gamma G_z z M_{xy} + i\gamma B_1(t) M_0 \qquad M_{xy} = M_x + i M_y \end{aligned}$$

First order linear differential equation. Easily solved.

 $\boldsymbol{\lambda}$ 

y

dN

dt

$$\frac{dM_{xy}}{dt} = -i\gamma G_z z M_{xy} + i\gamma B_1(t) M_0$$

Solving a first order linear differential equation:

$$M_{xy}(t,z) = i\gamma M_0 \int_0^t B_1(s) e^{-i\gamma G_z z \cdot (t-s)} ds$$
$$M_r(\tau,z) = iM_0 e^{-i\omega(z)\tau/2} \cdot \mathcal{FT}_{1D} \{\omega_1(t+\frac{\tau}{2})\} |_{t=-(\gamma/2\pi)G_z}$$

#### (See the note for complete derivation)

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$$M_{r}(\tau, z) = i M_{0} e^{-i\omega(z)\tau/2} \cdot \mathcal{FT}_{1D} \{ \omega_{1}(t + \frac{\tau}{2}) \} |_{f = -(\gamma/2\pi)G_{z}z}$$

To the Board



- For small tip angles, "the slice or frequency profile is well approximated by the Fourier transform of B1(t)"
- The approximation works surprisingly well even for flip angles up to 90°

## **Small Tip Approximation**

the excitation profile, within the small angle approximation, is just the Fourier transform of the pulse

remember that the Bloch equations are non-linear and thus cannot be expected to behave linearly

the approximation works surprisingly well even for flip angles up to 90°

### **Shaped Pulses**

 $30^{\circ}$  $90^{\circ}$ 0.6 0.5 0.8 0.4 0.6  $M_{v}$ Amplitude 0.3 0.4 Amplitude 0.2  $M_{\nu}$ 0.2  $M_{\rm x}$ 0.1  $M_{x}$ -0.2 -0.1-0.4-1 -0.8 -0.6 -0.4 -0.2n 0.2 0.4 0.6 0.8 -0.8 -0.6 -0.4 -0.2 0 0.2 0.4 0.8 0.6 Position Position

Pauly, J. J. Magn. Reson. 81 43-56 (1989)

small-angle approximation still works reasonably well for flip angles that aren't necessarily "small"

### **Truncation Artifacts**

in MRI we want pulses to be as short as possible to avoid relaxation effects

the sinc function is defined over all time which is impractical in any experiment

the sinc pulse needs to be truncated to be appropriate for clinical scans

### **Truncation Artifacts**

what happens when we truncate our pulses?



these deviations from the ideal are known as truncation artifacts

### **Truncation Artifacts**

#### alternative Pulse Shapes

gaussian

 $B_x(t) = A \exp\left[-a(t-\tau/2)^2\right]$ 

reduced side-lobes, but not as flat of a profile Window Functions

Hamming, Hanning, ...

## **Slice Rewinder**







# Slice Selective Excitation Example



slice select gradient rewinder eliminates the linear phase ramp





- Signal-to-Noise Ratio (SNR)
  - A fundamental measure of image quality

 $SNR \triangleq \frac{signal \ amplitude}{\sigma \ of \ noise}$ 

-  $SNR_{dB} = 20 \cdot log(SNR)$ 

#### Effect of Acquisition Time

- Simple 1D example (impulse in image space)
- N samples in k-space, each with amplitude A
- Noise variances add (independence)

$$SNR = \frac{\sum_{j=1}^{N} A}{\sqrt{\sum_{j=1}^{N} \sigma_n^2}} = \frac{NA}{\sqrt{N\sigma_n^2}} = \frac{\sqrt{NA}}{\sigma_n}$$

- Effect of Signal Averaging
  - Average separate measurements of the same kspace data samples (e.g., 2 measurements)
  - Signal amplitudes add

 $\Delta \Pi V E$ 

Noise variances also add (independence)

$$SNR_{2Ave} = \frac{\sum_{j=1}^{N} 2A}{\sqrt{\sum_{j=1}^{N} 2\sigma_n^2}} = \frac{2NA}{\sqrt{2N\sigma_n^2}} = \frac{\sqrt{2NA}}{\sigma_n}$$
$$SNR_{2Ave} = \sqrt{2} \cdot SNR$$

- Effect of Readout Time
  - Double readout duration  $T_{read}$
  - Typically, also double sampling interval  $\Delta t$  to maintain k-space sampling extent
  - $\Delta f \propto 1/(\Delta t)$ : halves the signal bandwidth  $\Delta f$

- Recall that 
$$\sigma_n^2 \propto \Delta f$$

$$SNR_{2 \cdot Tread} = \frac{NA}{\sqrt{N\sigma_n^2/2}} = \frac{\sqrt{2NA}}{\sigma_n}$$
$$SNR_{2 \cdot Tread} = \sqrt{2} \cdot SNR$$

- Summary of Acquisition Time Effects -  $SNR \propto \sqrt{N_{ave} \cdot T_{read}}$ -  $SNR \propto \sqrt{measurement time}$
- Effect of Spatial Resolution
  - $SNR \propto (\delta_x)(\delta_y)(\delta_z)$
- Other factors
  - $SNR \propto f(\rho, T_1, T_2, ...)$

To the Board

# **Gradient Echo Imaging**







- FID Decay due to
  - T2 decay
  - Spin
     dephasing
- Gradients accelerate spin dephasing
- Gradients can undo gradient induced spin dephasing



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 $T_2^*$  is signal loss from spin dephasing and  $T_2$ 

T2\*<T2 (always!)



# SE vs. GRE: B<sub>0</sub> Inhomogeneity

- Images acquired with a bad shim
  - Poor B<sub>0</sub> homogeneity (lots of off-resonance)





Spin Echo

**Gradient Echo** 

Images Courtesy of <u>http://chickscope.beckman.uiuc.edu/roosts/carl/</u> artifacts.html



- Related reading materials
  - Nishimura Chap 6 and 7

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