

## \* Magnetic fields

①

- $B_0$  large static field
- $B_1$  radio frequency field
- $G_{x,y,z}$  gradient fields

$B_0$  without  $B_0 \rightarrow$  spin oriented randomly

$$\sum \vec{\mu} = 0 = \vec{M}$$

↑

"net or bulk"  
magnetization

With  $B_0 \rightarrow$  2 things happen

a) Polarization

$$\begin{aligned} \sum \vec{\mu} \neq 0 &= \vec{M} \quad \text{alignment of spins} \\ &= M_z \hat{k} \end{aligned}$$

↑ along  $\hat{z}$  direction, "longitudinal"

In the presence of  $B_0$ , 2 energy states

$n_+$ , parallel

$n_-$ , antiparallel

\* Boltzmann distribution  $\frac{n_-}{n_+} = e^{-\Delta E/kT}$

(2)

b) Resonance

at equilibrium,  $\vec{M} \parallel \vec{B}$   
 If  $\vec{M} \times \vec{B}$ , precession will occur.

From classical mechanics,

$$\text{torque experienced by } \vec{M} \\ = \vec{M} \times \vec{B} = \frac{d}{dt} (\text{angular momentum}) = \frac{d}{dt} (\hbar \vec{I})$$

$$\vec{M} = \gamma \hbar \vec{I}$$

$$\Rightarrow \frac{d\vec{M}}{dt} = \vec{M} \times \gamma \vec{B}$$

↓ collection of spins

$$\frac{d\vec{M}}{dt} = \vec{M} \times \gamma \vec{B}, \quad \vec{B} = \vec{B}_0 = B_0 \hat{k}$$

$\Rightarrow \vec{M}$  will precess about the axis of  $\vec{B}$   
 at an angular frequency of  $\gamma |\vec{B}|$ .

\* frequency of precession

$$f = \frac{\gamma}{2\pi} B \text{ Hz ; Larmor frequency}$$

$$\text{For } ^1\text{H}, \quad \frac{\gamma}{2\pi} = \gamma^* = 42.575 \frac{\text{MHz}}{\text{T}}$$

(3)

$$\frac{d\vec{M}}{dt} = \vec{M} \times \gamma \vec{B}$$

$$\text{let } \vec{B} = B_0 \cdot \hat{k}$$

$$\frac{d\vec{M}}{dt} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ M_x & M_y & M_z \\ 0 & 0 & \gamma B_0 \end{vmatrix}$$

$$\frac{dM_x}{dt} = M_y \cdot \gamma B_0, \quad \frac{dM_y}{dt} = -M_x \cdot \gamma B_0, \quad \frac{dM_z}{dt} = 0$$

$$\begin{aligned} \frac{dM_x^2}{dt^2} &= \frac{dM_y}{dt} \cdot \gamma B_0, & \frac{dM_y^2}{dt^2} &= -\frac{dM_x}{dt} \cdot \gamma B_0 \\ &= -(\gamma B_0)^2 M_x & &= -(\gamma B_0)^2 M_y \end{aligned}$$

$$\text{Assume, } M_x(t) = A \cos(\gamma B_0 t) + B \sin(\gamma B_0 t)$$

$$\vec{M}^0 = \begin{bmatrix} M_x^0 \\ M_y^0 \\ M_z^0 \end{bmatrix}$$

$$M_x(t=0) = A = M_x^0$$

$$\begin{aligned} \frac{dM_x}{dt} &= -A \gamma B_0 \sin(\gamma B_0 t) + B \gamma B_0 \cos(\gamma B_0 t) \\ &= M_y \cdot \gamma B_0 \end{aligned}$$

$$M_y(t=0) = B = M_y^0$$

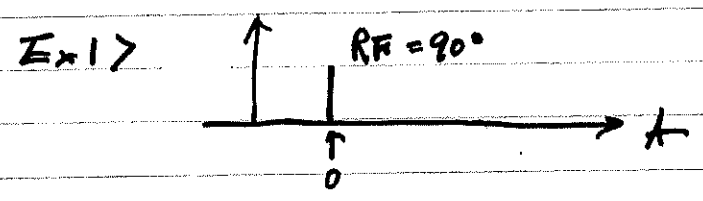
Finally,

$$\begin{aligned}
 M_x(t) &= M_x^0 \cos(\gamma B_0 t) + M_y^0 \sin(\gamma B_0 t) \\
 M_y(t) &= -M_x^0 \sin(\gamma B_0 t) + M_y^0 \cos(\gamma B_0 t) \\
 M_z(t) &= M_z^0
 \end{aligned}$$

or,

$$\begin{bmatrix} M_x(t) \\ M_y(t) \\ M_z(t) \end{bmatrix} = \underbrace{\begin{bmatrix} \cos(\gamma B_0 t) & \sin(\gamma B_0 t) & 0 \\ -\sin(\gamma B_0 t) & \cos(\gamma B_0 t) & 0 \\ 0 & 0 & 1 \end{bmatrix}}_{R_z(\gamma B_0 t)} \begin{bmatrix} M_x^0 \\ M_y^0 \\ M_z^0 \end{bmatrix}$$

$$\begin{aligned}
 \omega &= \gamma B_0 \\
 \vec{\omega} &= \gamma \vec{B}_0 = \gamma B_0 \hat{k}
 \end{aligned}$$



$$\vec{M}(0_-) = \begin{bmatrix} 0 \\ 0 \\ M_z^0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ M_0 \end{bmatrix}$$

$$\vec{M}(0_+) = \begin{bmatrix} M_0 \\ 0 \\ 0 \end{bmatrix} \leftarrow \text{Immediately after a } 90^\circ \text{ RF pulse}$$

$$\vec{M}(t) = R_z(\gamma B_0 t) \begin{bmatrix} M_0 \\ 0 \\ 0 \end{bmatrix} \leftarrow \vec{M}(t) \text{ after the } 90^\circ \text{ RF pulse}$$

$$\begin{aligned}
 M_x(t) &= M_0 \cos(\gamma B_0 t) \\
 M_y(t) &= -M_0 \sin(\gamma B_0 t) \\
 M_z(t) &= 0
 \end{aligned} \leftarrow \text{free precession}$$